

A Appendix: Analytic Solution of the Minimal Model

It helps to introduce a new parameterization of time:

$$u = \log \phi(t) = \log \left(1 + \frac{(R+1)t}{N_0} \right) \quad (20)$$

with associated derivative:

$$\frac{\partial}{\partial t} = \frac{R+1}{N_0 e^u} \frac{\partial}{\partial u} \quad (21)$$

With this definition, $u = 0$ corresponds to $t = 0$.

This new variable helps rid the differential equations (2) of explicit time dependence:

$$\begin{aligned} \frac{\partial F(m, u)}{\partial u} + \frac{mF(m, u)}{R+1} &= \frac{(m-1)F(m-1, u)}{R+1} \quad (m > 1) \\ \frac{\partial F(1, u)}{\partial u} + \frac{F(1, u)}{R+1} &= \frac{N_0 R}{R+1} e^u \end{aligned} \quad (22)$$

Note that the equation the special bin $F(1, u)$ does not depend on any other $F(m, u)$, so it can be solved separately. Once it is known, the solution for any other m can be found by successive integration:

$$F(m+1, u) = \exp \left(-\frac{m+1}{R+1} u \right) \int_0^u \frac{dv}{R+1} mF(m, v) \exp \left(\frac{m+1}{R+1} v \right) \quad \text{for } m+1 = 2, 3, \dots \quad (23)$$

The solution for m serves as a “source” for $m+1$. The relation (23) follows by multiplying both sides of (2) by $\exp \left(\frac{m+1}{R+1} u \right)$ and integrating. Note that this solution ensures that $F(m > 1, t = 0) = 0$, so the initial conditions are automatically satisfied. Our method of solving the differential equation is elementary and standard, see (Simmons, 1994) for more details.

The solution for $m = 1$ can be found in the same way:

$$\frac{\partial}{\partial u} \left[\exp \left(\frac{u}{R+1} \right) F(1, u) \right] = \exp \left(\frac{u}{R+1} + u \right) \frac{N_0 R}{R+1} \quad (24)$$

$$F(1, u) = N_0 \exp \left(-\frac{u}{1+r} \right) + \frac{N_0 R}{R+2} \left[\exp u - \exp \left(-\frac{u}{1+r} \right) \right] \quad (25)$$

The full solution follows by successive application of (23). There are two types of integrals that come up:

$$\begin{aligned} \exp \left(-\frac{m+1}{R+1} u \right) \int_0^u \frac{dv}{R+1} [m \exp v] \exp \left(\frac{m+1}{R+1} v \right) \\ = \frac{m}{R+2+m} \left[\exp u - \exp \left(-\frac{m+1}{R+1} u \right) \right] \end{aligned} \quad (26)$$

$$\begin{aligned} \exp \left(-\frac{m+n+1}{R+1} u \right) \int_0^u \frac{dv}{R+1} (m+n) \\ \left[\exp \left(-\frac{mv}{R+1} \right) \left(1 - \exp \left(-\frac{v}{R+1} \right) \right)^n \right] \exp \left(\frac{m+n+1}{R+1} v \right) \\ = \frac{m+n}{n+1} \left[\exp \left(-\frac{mu}{R+1} \right) \left(1 - \exp \left(-\frac{u}{R+1} \right) \right)^{n+1} \right] \end{aligned} \quad (27)$$

The coefficients that emerge from these integrations define the recursion relations for A_m and β_n^m :

$$A_{m+1} = \frac{m}{R+2+m} A_m \quad (28)$$

$$\beta_{n+1}^m = \frac{m+n}{n+1} \beta_n^m \quad (29)$$

The full solution to (2), taking into account the initial conditions, is given by:

$$F(m, t) = N_0 \phi^{-\frac{1}{1+R}} \left(1 - \phi^{-\frac{1}{1+R}}\right)^{m-1} + A_m (\phi - \phi^{-\frac{m}{1+R}}) - \sum_{i=1}^{m-1} A_i \beta_{m-i}^i \phi^{-\frac{i}{1+R}} \left(1 - \phi^{-\frac{1}{1+R}}\right)^{m-i}$$

$$A_m = \frac{RN_0}{R+2} \prod_{i=1}^{m-1} \frac{i}{R+2+i} = RN_0 \frac{\Gamma(m)\Gamma(R+2)}{\Gamma(R+2+m)}$$

$$\beta_n^m = \prod_{k=1}^n \frac{m+k-1}{k} = \frac{(m+n-1)!}{(m-1)!n!} \quad (30)$$

with the understanding that an empty product is unity, i.e. $\prod_{i=1}^0 f(i) = 1$.

Note that the coefficients A_m and β_n^m do not depend on time, and furthermore β_n^m has no dependence on R or N_0 . The product of coefficients $A_i \beta_{m-i}^i$ can be simplified:

$$A_i \beta_{m-i}^i = \frac{(m-1)!}{(m-i)! \prod_{j=1}^{i-1} (R+2+j)} \quad (31)$$

but it will be useful keep these these coefficients separate when considering the solution for more general initial conditions. Note that we use the standard definition for the gamma function $\Gamma(x)$; see Appendix H.

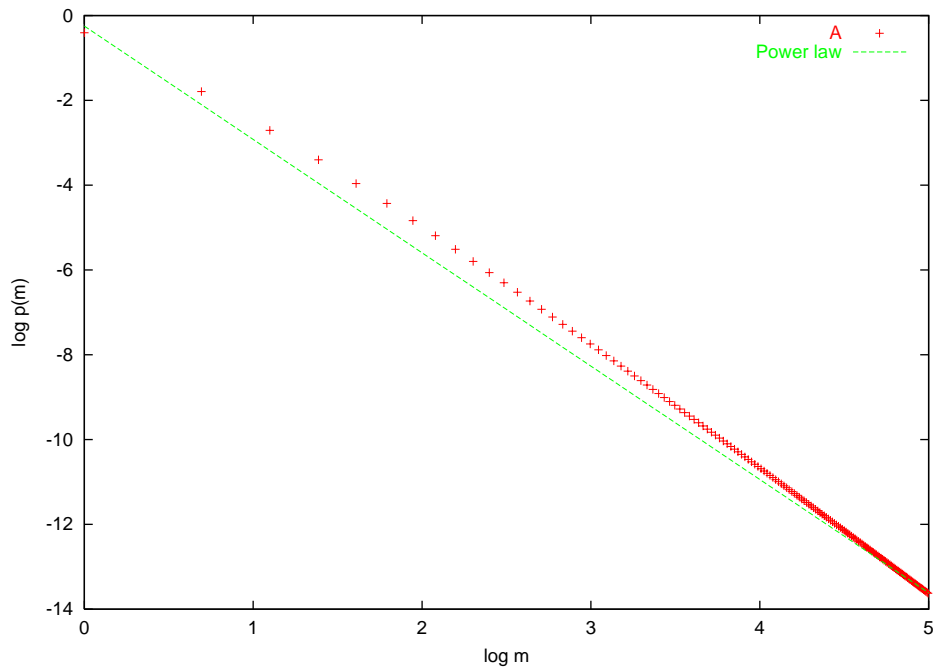


Figure 6: The normalized A_m coefficients (points) and a power-law fit (line), shown as a log-log plot as a function of size m , for $R = 1$.